

Can all the observed neutron stars contain color-superconducting quark cores?

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Credits: **A. Ayriyan & D. Blaschke**

Compact Stars in QCD Phase Diagram 2026

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PHYSICAL REVIEW D **111**, 034004 (2025)

Asymptotically conformal color-flavor-locked quark matter within a nonlocal chiral quark model

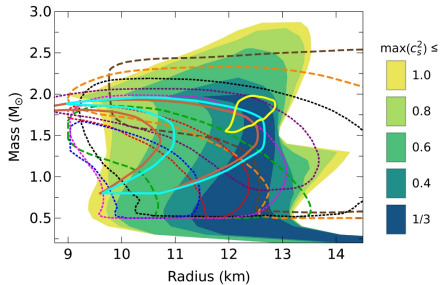
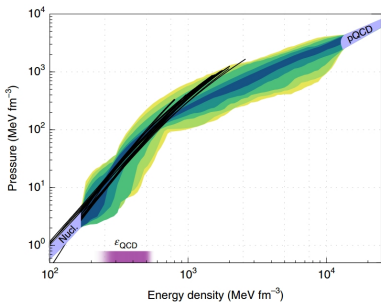
Oleksii Ivanytskyi ^{*}

arxiv:2508.02554 [nucl.th], accepted to Physical Review D

Bayesian inference favors quark matter in neutron star interiors

Alexander Ayriyan ^{1,2,*} Oleksii Ivanytskyi ^{1,3,4,5,†} and David Blaschke ^{1,4,5,‡}

Quark matter in neutrons stars? pQCD vs $2M_{\odot}$



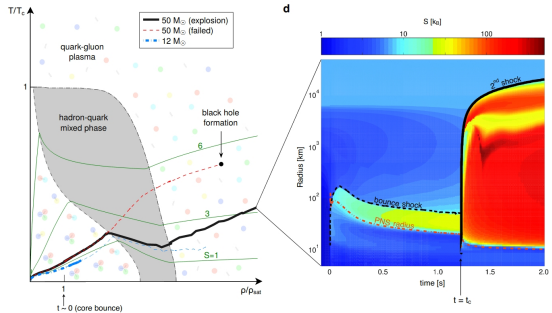
E. Annala, T. Gorda, A. Kurkela, J. Nättilä, A. Vuorinen, *Nature Physics* 16, 907 (2020)

Existence of parameterization consistent with pQCD and $2M_{\odot}$

Can NSs accommodate quark cores?

Quark matter in neutrons stars? Supernova explosions

- $2M_{\odot}$ stars formation? (accretion is too slow)
- Supernovae with progenitor mass $\sim 50 M_{\odot}$
- Quark-hadron transition stabilizes collapse



T. Fischer et al., Nature Astronomy 2, 980–986 (2018)

Table 1 | Summary of the supernova simulation results with hadron-quark phase transition

M_{ZAMS} (M_{\odot})	t_{onset} (s)	t_{collapse} (s)	ρ_{collapse} (ρ_{sat})	T_{collapse} (MeV)	$M_{\text{PNS,collapse}}^a$ (M_{\odot})	t_{final} (s)	ρ_{final} (ρ_{sat})	T_{final} (MeV)	$M_{\text{PNS,final}}^a$ (M_{\odot})	E_{expl} (10^{51} erg)
12 ¹²	3.251	3.489	2.49	28	1.727	3.598	5.5	17	1.732	0.1
18 ¹²	1.465	1.518	2.53	27	1.958	1.575	5.9	18	1.964	1.6
25 ¹⁴	0.905	0.976	2.40	31	2.163	0.983	9.6	19	2.171 ^b	-
50 ¹	1.110	1.215	2.37	32	2.105	1.224	5.8	31	2.092	2.3

Deconfinement is a supernova engine for blue supergiants

Quark matter in neutrons stars? Hyperons

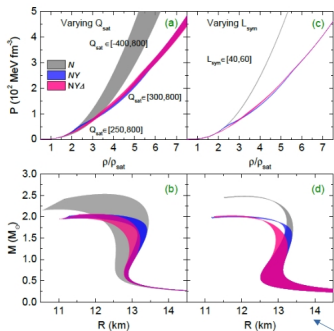


FIG. 4. EoS models and MR relations for N , NY , and NY_{Δ} compositions of stellar matter. The bands are generated by varying the parameters Q_{sat} [MeV] (a, b) and L_{sym} [MeV] (c, d). The ranges of Q_{sat} and L_{sym} allowed by χ EFT and maximum mass constraints are indicated in the figures.

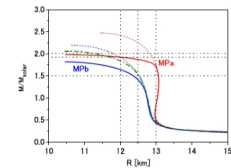


FIG. 7. Neutron-star masses as a function of the radius R . Solid (dashed) curves are with (without) hyperon (Λ and Σ^-) mixing for ESC+MPs and ESC+MPb. The dot-dashed curve for MPb is with Λ mixing only. Also see the caption of Fig. 3.

Yamamoto et al., Phys.Rev.C 96 (2017) 06580;
arXiv:1708.06163 [nucl-th]
Yamamoto et al., Eur. Phys. J. A 52 (2016) 19;
arXiv:1510.06099 [nucl-th]
Ji & Sedrakian, Phys. Rev. C 100 (2019) 015809;
arXiv:1903.06057 [astro-ph.HE]

Examples for realistic hadronic EoS which suggest a Berlin Wall is inferior to the line $M = 2.0 M_{\odot}$

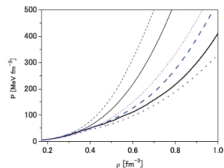


Fig. 8. Pressure P as a function of baryon density ρ . Thick (thin) curves are with (without) hyperon mixing. Solid, dashed and dotted curves are for MPs, MPb and MPb.

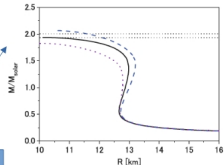


Fig. 9. Neutron-star masses as a function of the radius R . Solid, dashed and dotted curves are for MPs, MPb and MPb. Two dotted lines show the observed mass (1.97 ± 0.04) M_{\odot} of J1614-2230.

Hyperons soften EoS. Can quarks stiffen it?

Model to probe quark matter neutron stars

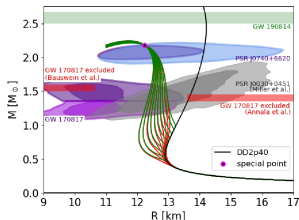
- Quark degrees of freedom
- Asymptotically conformal behavior in agreement with pQCD
- Consistency with observational constraints on properties of neutron stars

$N_f = 3$ nonlocal NJL model with vector repulsion and diquark pairing

- Realistic hadronic EoS: **DD2npY-T**
- Phase transition within Maxwell construction (not principal for modeling NSs)

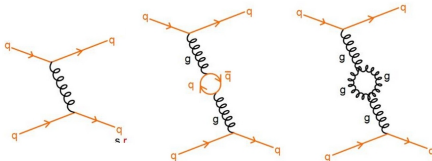
M. Shahrhaf et al., PRD (2022)

M. Cierniak, D. Blaschke, Astron.Nachr. 342 (2021)



Nonlocal character of quark interactions

- QCD interactions

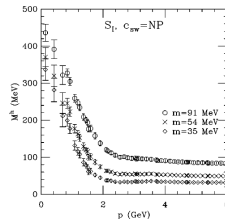


- Indications from the lattice
 p -dependent renormalization of quark propagator
 (p -dependent quark mass)



space-time nonlocality of quark interactions

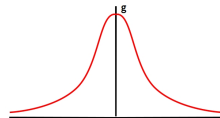
J. Skullerud, D. B. Leinweber, and A. G. Williams, *Phys. Rev. D* 64 , 074508 (2001)



- Quark interactions in effective approaches (separable approximation)

$$\left(\bar{q}_x \hat{\Gamma} q_x\right)^2 \rightarrow \left(\int dz \tilde{g}_z \bar{q}_{x+z/2} \hat{\Gamma} q_{x-z/2}\right)^2$$

\tilde{g}_z – space-time dependent formfactor



Nonlocal NJL model for three-flavor quark matter

- Lagrangian

$$\mathcal{L} = \bar{q}(i\cancel{\partial} - m + \mu\gamma_0)q + \overbrace{G_S \sum_{a=0,8} s_a s_a}^{\text{chiral dynamics}} - \overbrace{G_V j_{\mu} j^{\mu}}^{\text{vector repulsion}} + \overbrace{3G_D \sum_{a,b=2,5,7} d_{ab}^+ d_{ab}}^{\text{color superconductivity}}$$

$$\int dz \tilde{g}(z) \bar{q}\left(x + \frac{z}{2}\right) \hat{\Gamma} q\left(x - \frac{z}{2}\right) = \begin{cases} s_a(x), & \hat{\Gamma} = \tau_a \\ j_{\mu}(x), & \hat{\Gamma} = \gamma_{\mu} \\ d_{ab}(x), & \hat{\Gamma} = i\gamma_5 \tau_a \lambda_b \hat{C} \end{cases}$$

- Thermodynamic potential

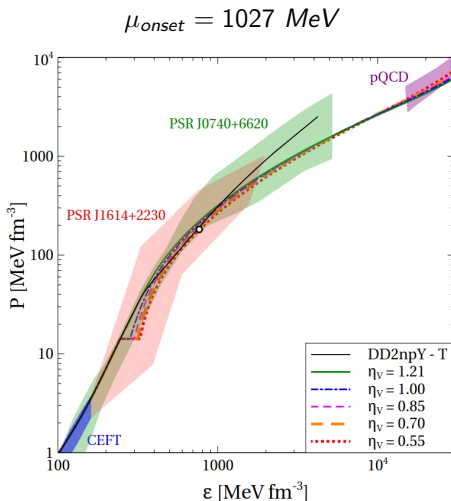
$$\Omega = - \sum_{j,a=\pm} d_j \int_{\mathbf{k}} \left[\frac{1}{2} - f_{j\mathbf{k}}^a \right] \epsilon_{j\mathbf{k}}^a + \frac{\sigma^2}{4G_S} - \frac{\omega^2}{4G_V} + \frac{\Delta^2}{4G_D}$$

$$\epsilon_{j\mathbf{k}}^a = \text{sign}(\epsilon_{\mathbf{k}}^a) \sqrt{\epsilon_{\mathbf{k}}^{a2} + \Delta_{j\mathbf{k}}^2}, \quad \epsilon_{\mathbf{k}}^{\pm} = \sqrt{\mathbf{k}^2 + M_{\mathbf{k}}^2} \mp \mu \mp \omega g_{\mathbf{k}}, \quad M_{\mathbf{k}} = m + \sigma g_{\mathbf{k}}$$

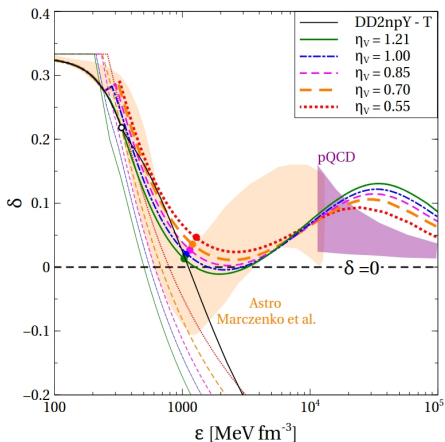
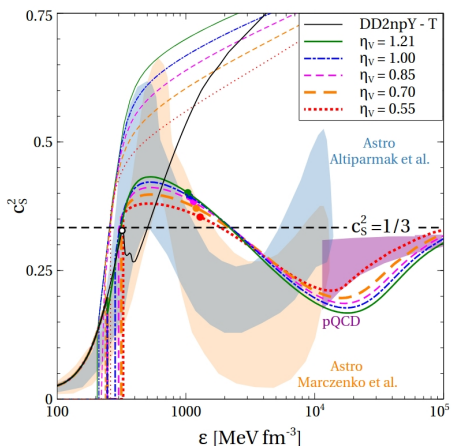
$$\frac{\partial \Omega}{\partial \sigma} = \frac{\partial \Omega}{\partial \omega} = \frac{\partial \Omega}{\partial \Delta} = 0$$

EoS of neutron stars

- Charge neutrality
- β -equilibrium
- Hadron EoS: DD2npY-T
- Maxwell construction



Speed of sound & dimensionless conformality measure



c_s^2 and δ almost simultaneously attain the conformal values
Pseudoconformality at nonperturbative densities (not discussed here)

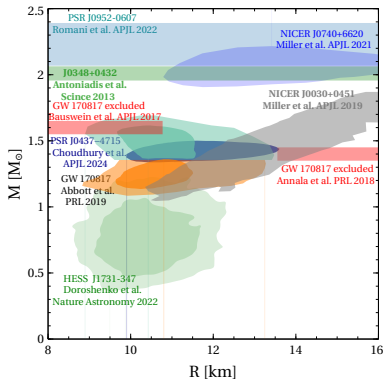
Model-informed Bayesian analysis: data

• Basic constraints

- 1 PSR J0348+0432 by Antoniadis et al.
- 2 PSR J0030+0451 by Miller et al.
- 3 PSR J0740+6620 by Dittmann et al.
- 4 PSR J0437-4715 by Choudhury et al.
- 5 GW170817 by Abbott et al.

• Additional constraints

- 1 PSR J0952-0607 by Romani et al.
- 2 HESS J1731-347 by Doroshenko et al.

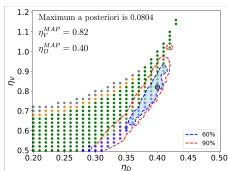
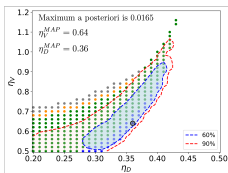


Model-informed Bayesian analysis: results

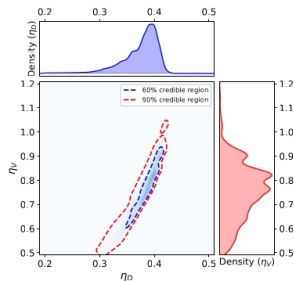
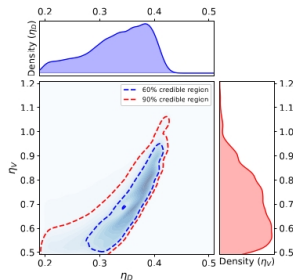
- $\eta_V = [0.5, 1.5], \eta_D = [0.2, 0.5]$

- $51 \times 31 = 1581$ quark EoSs

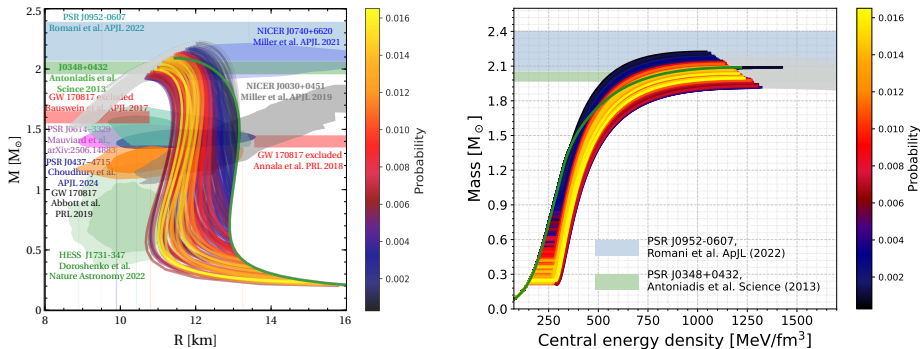
- 259 hybrid EoSs



- $\eta_V \times \eta_D - s$ space



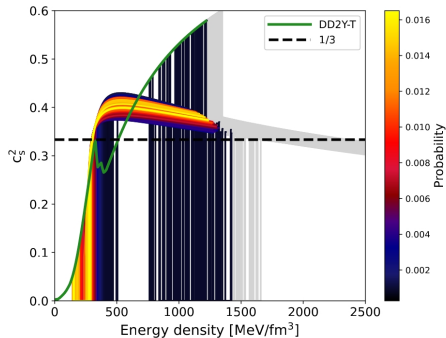
Neutron stars with CFL cores



A. Ayriyan, O. Ivanytskyi, D. Blaschke, 2509.02554 [nucl-th], accepted to PRD

Onset of quark matter below $\epsilon \simeq 300 \text{ MeV}/\text{fm}^3$ and $1M_{\odot}$
All the observed neutron stars have quark cores

Speed of sound of the neutron star matter



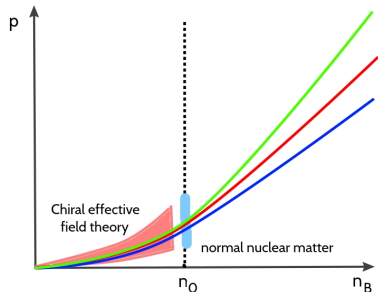
A. Ayriyan, O. Ivanytskyi, D. Blaschke, 2509.02554 [nucl-th], accepted to PRD

Speed of sound peaks at $\varepsilon \simeq 500 \text{ MeV/fm}^3$ and $c_s^2 \simeq 1/3$
Neutron star matter is far from conformality

The role of hadronic EoS

• Constraints

- 1 Chiral effective field theory
- 2 Normal nuclear matter
- 3 Symmetry energy from PREX
- 4 ...



**All realistic nuclear EoSs should be very similar at $1 - 2 n_0$.
Thus, make no difference if $n^{\text{onset}} < 2n_0$**

Conclusions

- Hybrid stars are 1-2 o.o.m. more preferred over the hadronic stars
- Most probable onset of quark matter

$$n_{\text{onset}} \lesssim 2n_0, \quad M \lesssim 1M_{\odot}$$

All neutron stars have quark cores

- My proposal of phase diagram

Quark matter should be included to simulations of supernovae, NS cooling and mergers

